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Measurement of the direct CP -violating parameter A_{CP} in the decay $D^+ \rightarrow K^- \pi^+ \pi^+$

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We measure the direct CP -violating parameter A_{CP} for the decay of the charged charm meson, $D^+ \rightarrow K^- \pi^+ \pi^+$ (and charge conjugate), using the full 10.4 fb^{-1} sample of $p\bar{p}$ collisions at $\sqrt{s} = 1.96 \text{ TeV}$ collected by the D0 detector at the Fermilab Tevatron collider. We extract the raw reconstructed charge asymmetry by fitting the invariant mass distributions for the sum and difference of charge-specific samples. This quantity is then corrected for detector-related asymmetries using data-driven methods and for possible physics asymmetries (from $B \rightarrow D$ processes) using input from Monte Carlo simulation. We measure $A_{CP} = [-0.16 \pm 0.15(\text{stat}) \pm 0.09(\text{syst})]\%$, which is consistent with zero, as expected from the standard model prediction of CP conservation, and is the most precise measurement of this quantity to date.

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The violation of CP symmetry in the fundamental interactions of particle physics is required to explain the matter dominance of the Universe [1–3]. The standard model (SM) describes CP violation in the quark sector through the presence of a single irreducible phase in the Cabibbo-Kobayashi-Maskawa (CKM) matrix. This matrix dictates the strength of flavor transitions through the weak interaction. All experimental observations to date are

consistent with a single phase [4], with the exception of a small number of discrepancies at the $\approx 3\sigma$ level, most notably the anomalously large same-charge dimuon asymmetry measurement from the D0 experiment [5]. However, the degree of CP violation in the SM is insufficient to explain the cosmological matter dominance [6]. It is therefore important to continue searching for sources of CP violation beyond those predicted by the SM.

Decays of heavy-flavor hadrons provide a natural testing ground for these searches. In particular, decays proceeding through box or penguin diagrams are highly sensitive to possible CP violation contributions from processes beyond the SM induced by additional particles in the loops. However, due to the difficulty in simultaneously extracting production, detection and physics asymmetries, these searches for anomalous CP violation typically measure the difference in charge asymmetries between the channel of interest and a Cabibbo-favored reference channel, which is then assumed to be CP symmetric [7–11]. Performing high-precision measurements of CP violation parameters in these Cabibbo-favored decays is therefore crucial in order to establish an experimental basis for these assumptions, thus reducing dependence on theoretical predictions. The data set collected by the D0 experiment at the Tevatron $p\bar{p}$ collider is uniquely suited to perform such measurements, having a CP -symmetric initial state and almost equal beam exposure in all four combinations of solenoid and toroid magnet polarities.

In this paper, we describe the measurement of the direct CP violation parameter in the Cabibbo-favored decay $D^+ \rightarrow K^- \pi^+ \pi^+$ (charge conjugate states are implied throughout this paper), defined as

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$$A_{CP}(D^+ \rightarrow K^- \pi^+ \pi^+) = \frac{\Gamma(D^+ \rightarrow K^- \pi^+ \pi^+) - \Gamma(D^- \rightarrow K^+ \pi^- \pi^-)}{\Gamma(D^+ \rightarrow K^- \pi^+ \pi^+) + \Gamma(D^- \rightarrow K^+ \pi^- \pi^-)} \quad (1)$$

and hereafter denoted A_{CP} . Currently this parameter has only been measured by the CLEO Collaboration [12]: $A_{CP} = [-0.3 \pm 0.2(\text{stat}) \pm 0.4(\text{syst})]\%$. We use the complete sample of $p\bar{p}$ collisions generated by the Tevatron accelerator at $\sqrt{s} = 1.96$ TeV and collected by the D0 detector using a suite of muon triggers. This corresponds to approximately 10.4 fb^{-1} of integrated luminosity.

CP violation can only occur if there is interference between two amplitudes with different strong and weak phases. For the decay mode being investigated, this requirement is not satisfied, with two tree-level amplitudes both proportional to the product of CKM matrix elements $V_{cs}^* V_{ud}$ and no contribution from Cabibbo-suppressed diagrams. The SM therefore predicts negligible CP violation with respect to the experimental uncertainties. Any significant deviation of A_{CP} from zero would thus constitute evidence for new physics contributions [13].

Experimentally, the CP asymmetry parameter is determined by measuring a raw charge asymmetry (A) and applying corrections to account for differences in the detection of the final-state particles (A_{det}) and in the production rates of D^+ and D^- mesons (A_{phys}), i.e., neglecting terms of second-order or higher in the asymmetries,

$$A_{CP} = A - A_{\text{det}} - A_{\text{phys}}. \quad (2)$$

The raw quantity A is the asymmetry in the number of D^+ versus D^- mesons reconstructed in the described decay mode and passing all selection requirements. It is extracted by simultaneously fitting the $M(K\pi\pi)$ invariant mass distributions for the sum of all candidates and for the difference $N(D^+) - N(D^-)$. The detector asymmetry A_{det} accounts for differences in the reconstruction efficiency for positive and negative kaons, pions, and muons and is determined using methods based on data in dedicated independent channels. The physics asymmetry A_{phys} accounts for possible charge-asymmetric production of D mesons arising through the decay of B hadrons. For each possible source, the contribution to A_{phys} is the product of the relevant CP asymmetry (taken from the world average of experimental results) and the fraction of D mesons arising from this source (determined from simulation). We assume negligible CP violation in the decays of B mesons into final states containing D^\pm . In practice the value of A_{phys} is small compared to the precision of the final measurement, while the detector correction is significant. For simplicity, we use D to collectively denote D^\pm mesons throughout this paper. In cases where distinguishing the charge is important we explicitly include it.

The D0 detector is described in detail elsewhere [14,15]. The most important components for this analysis are the central tracking detector, the muon system, and the

magnets. The central tracking system comprises a silicon microstrip tracker (SMT) closest to the beam pipe, surrounded by a central fiber tracker (CFT), with the entire system located within a 1.9 T solenoidal field. The SMT (CFT) has polar acceptance $|\eta| < 3$ ($|\eta| < 2.5$), where the pseudorapidity is defined as $\eta = -\ln[\tan(\theta/2)]$, and θ is the polar angle with respect to the positive z axis along the proton beam direction. The muon system (covering $|\eta| < 2$) comprises a layer of tracking detectors and scintillation trigger counters in front of 1.8 T toroid magnets, followed by two similar layers after the toroids. The polarities of both the solenoid and toroid magnets were regularly reversed approximately every two weeks during data collection to give near equal exposure in all four configurations. The magnet reversal ensures that the main detector asymmetries cancel to first order by symmetrizing the detector acceptance for positive and negative particles. The residual deviations from equal exposure (typically less than 5%) are removed by weighting events to force equal contributions from all four polarity configurations.

In the absence of a dedicated trigger for hadronic decays of heavy-flavor hadrons, we use a suite of single muon and dimuon triggers to select the data sample, along with an off-line single muon filter. Events that exclusively satisfy triggers using track impact parameter information are removed to avoid lifetime biases which influence the D meson parentage, and which are challenging to model in simulation. The muon trigger and off-line requirements can bias the composition of the data in favor of semileptonic decays of charm and bottom hadrons. In particular, the fraction of D mesons arising from semileptonic decays of B mesons will be enhanced. These requirements must be taken into account when determining both detector and physics asymmetry corrections. To facilitate this process, the analysis places particular requirements on the muon quality and kinematic variables, to match those used when determining kaon, pion, and muon reconstruction asymmetries. The muon must produce hits in the muon tracking layers both inside and outside the toroid and must be spatially matched to a central track with total momentum $p(\mu) > 3 \text{ GeV}/c$ and transverse momentum $p_T(\mu) > 2 \text{ GeV}/c$. The selected muon is not used in the subsequent reconstruction of D meson signal candidates. In particular, no further requirements are imposed which use the muon information (for example, charge, or spatial origin with respect to the D meson candidate). For events with more than one muon (around 9%) the one with highest p_T is used consistently when determining the associated background asymmetries.

For events passing the muon selection, D candidates are reconstructed from all possible three-track combinations that have total charge $q = \pm 1$ and that are consistent with arising from a common vertex. The three tracks must satisfy quality requirements and each track must have $p_T > 0.7 \text{ GeV}/c$. The two like-charge tracks are assigned

the charged pion mass, and the third track is assigned the charged kaon mass [4]. The resulting invariant mass of the D candidate must lie within $1.65 < M(K\pi\pi) < 2.05$ GeV/ c^2 , and the momentum and displacement vectors of the reconstructed D meson must point in the same hemisphere. Additionally, the transverse decay length of the D candidate must exceed 3 times its uncertainty, $L_{xy}(D)/\sigma[L_{xy}(D)] > 3$. The transverse decay length is defined as the displacement between the $p\bar{p}$ primary interaction vertex and the reconstructed D meson decay vertex, projected onto the plane perpendicular to the beam direction.

The final selection of events uses a log-likelihood ratio (LLR) method to combine 12 individual variables into a single multivariate discriminant, using a similar approach to that described in Ref. [10]. The input variables are as follows: the transverse momenta of the three final-state hadrons and their track isolations, the transverse decay length of the D meson L_{xy} and its significance $L_{xy}(D)/\sigma[L_{xy}(D)]$, the χ^2 of the vertex fit of the three tracks, the angular separations of the kaon and lowest- p_T pion and of the two pions, and the cosine of the angle between the momentum and displacement vectors of the D meson candidate. The angular separation of two tracks is defined as $\Delta R = \sqrt{\Delta\phi^2 + \Delta\eta^2}$, where $\Delta\phi$ and $\Delta\eta$ are the track separations in the azimuthal angle and pseudorapidity, respectively. The track isolation I is the momentum of a particle divided by the sum of the momenta of all tracks contained in a cone of size $\Delta R < 0.5$ around the particle. Tracks corresponding to the other two final-state particles for this candidate are excluded from the sum.

Distributions observed in background-dominated data are used to derive the likelihood functions for background-like events, which are subsequently used in the LLR discriminant. These distributions are populated using 1% of the data, chosen by randomly sampling the D candidates following all requirements except for the LLR. The signal contamination in this sample is small enough (around 0.4%) that it does not affect the performance of the discriminant. The signal distributions are modeled using Monte Carlo (MC) simulation of inclusive $D^\pm \rightarrow K^\mp \pi^\pm \pi^\pm$ events, without any constraints on their origin. The final requirement on the LLR output is chosen to maximize the signal significance in the 1% random data sample (scaling up to extrapolate to the full sample). Ensemble studies confirm that this corresponds to the minimum statistical uncertainty on the final asymmetry measurement.

For all simulated samples, events are generated using PYTHIA version 6.409 [16] interfaced with EVTGEN [17] to model the decays of particles containing b and c quarks. The generation model includes all quark flavors, ensuring that charm and bottom quarks from gluon splitting are properly included in the final sample. Generated events are processed by a GEANT-based detector simulation [18], overlaid with data from randomly collected bunch

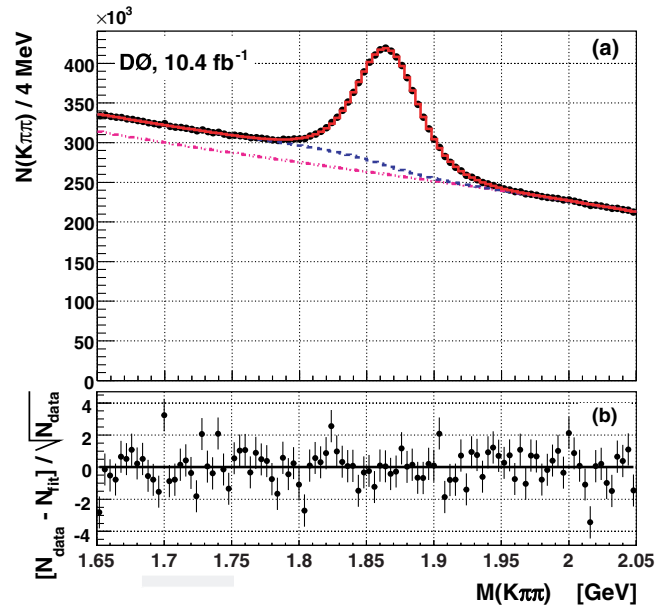


FIG. 1 (color online). (a) Invariant mass distribution $M(K\pi\pi)$ after all selections have been applied (data markers). Also shown is the result of the fit to the data, as described in the text (solid line). To illustrate the contributions of the three separate components, the total background (dashed line) and polynomial function (dot-dashed line) are shown separately. (b) Fit residuals $[N_{\text{data}} - N_{\text{fit}}]/\sqrt{N_{\text{data}}}$, demonstrating the agreement between the data and the fit model.

crossings to simulate pileup from multiple interactions, and reconstructed using the same software as used for data.

The $M(K\pi\pi)$ distribution of candidates passing all selection requirements is shown in Fig. 1, along with the results of a fit to the data (described later). A total of approximately 31×10^6 candidates is found, of which $N(D^\pm) = 2270224 \pm 7406$ are assigned as D^\pm signal in the fit. The effective statistical loss caused by the magnet polarity weighting, included in this number, is 3.2%. For around 10% of events we find multiple D^\pm candidates, all of which are accepted and treated independently when determining the raw and background asymmetries.

The raw asymmetry A is extracted through a simultaneous binned minimum- χ^2 fit of the sum distribution (in Fig. 1) and the difference distribution $[N(D^+) - N(D^-)]$ (in Fig. 2). The method is the same as described in Ref. [10], with the only difference being a slight simplification of the combinatorial background model, enabled by the updated event selection criteria. The fit includes three components, each having the same shape in the sum and difference distributions, with only their relative normalizations differing in the two cases. The D signal is parametrized by two Gaussian functions constrained to have the same mean value, to model the effect of the detector mass resolution. A hyperbolic tangent function is used to model the effect of a range of multibody physics backgrounds, including both partially reconstructed decays

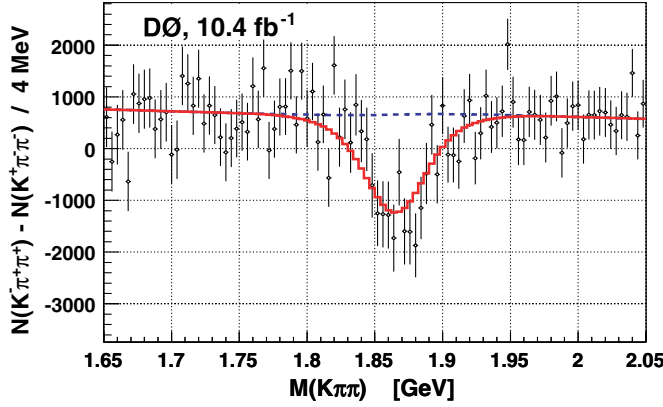


FIG. 2 (color online). Invariant mass distribution $M(K\pi\pi)$ for the difference $N(D^+) - N(D^-)$ (data markers). Also shown are the result of the fit to the data (solid line) and the overall background contribution (dashed line).

of $D^{(*)}$ mesons, and reflections where the final-state hadrons are assigned the wrong mass. The main contributions are from D^+ decays to $K^-\pi^+\pi^+\pi^0$, $\pi^-\pi^+\pi^+\pi^0$, and $K^-K^+\pi^+$; D_s^+ decays to $K^+K^-\pi^+$; \bar{D}^0 decays to four charged hadrons; and decays of $D^{*+} \rightarrow D^0\pi^+$, with $D^0 \rightarrow K^-\pi^+\pi^0$, where in all cases the π^0 is not reconstructed. The hyperbolic tangent parametrization is chosen based on studies of decay-specific and inclusive simulated samples and is the same as used in Ref. [10]. The inflection point is fixed for the nominal fit, based on simulation, but is allowed to vary when assigning a systematic uncertainty to the choice of fitting model. The steepness of the slope is constrained based on the resolution of the Gaussian peak in data [10], which is also well motivated by simulation. Finally, the smooth combinatorial background is modeled by a polynomial with constant, linear, and cubic terms. The quadratic term is excluded since it does not improve the goodness of fit. For the fit to the difference distribution, the relative contributions of the three components are quantified through asymmetry parameters, including the raw asymmetry A for the signal and corresponding asymmetries A_{multi} and A_{comb} for the multibody and combinatorial components, respectively. Hence the models used to fit the sum (F_{sum}) and difference (F_{diff}) distributions can be summarized as

$$\begin{aligned} F_{\text{sum}} &= F_{\text{sig}} + F_{\text{comb}} + F_{\text{multi}}, \\ F_{\text{diff}} &= A \cdot F_{\text{sig}} + A_{\text{comb}} \cdot F_{\text{comb}} + A_{\text{multi}} \cdot F_{\text{multi}}, \end{aligned} \quad (3)$$

where, for instance, F_{sig} is the function used to model the signal component.

The total number of candidates and the difference between the positive and negative candidate counts are used as constraints to reduce the number of free parameters by two, giving improved fit stability. The final fit has ten free parameters, six for the signal [signal yield $N(D)$,

invariant mass $M(D)$, the widths of the two Gaussian functions, the fraction of signal in the wider Gaussian, and the raw asymmetry] and four for the background (fraction of background in multibody component, first- and third-order polynomial coefficients, and A_{multi}). The final two variables, A_{comb} and the constant term in the polynomial function, are completely defined by the set of ten free parameters and the two external constraints.

The corresponding distribution and fit for the difference $[N(D^+) - N(D^-)]$ is shown in Fig. 2. A significant negative raw asymmetry is observed, $A = (-1.28 \pm 0.15)\%$, consistent with the value expected from known detector asymmetries. The two background asymmetries are $A_{\text{multi}} = (-0.41 \pm 0.60)\%$ and $A_{\text{comb}} = (+0.27 \pm 0.04)\%$. The main source of charge asymmetry in both background components is the kaon reconstruction asymmetry, which is around $+1.1\%$ and is described later. The sign and magnitude of both A_{multi} and A_{comb} are consistent with expectations from this kaon asymmetry alone. The main processes contributing to the multibody component, and including a single charged kaon in the final state, are from the Cabibbo-favored transition $c \rightarrow s$. This results in a negative correlation between the kaon and D charge, so we expect A_{multi} to be negative, with a magnitude somewhat less than 1.1% due to dilution from processes without a single charged kaon. In contrast, the combinatorial background component models the contribution of random three-track combinations: the kaon asymmetry leads to an overall excess of positive tracks, and so A_{comb} is expected to be positive, with a magnitude driven by the relative abundance of kaons in the track sample. The full fit to both distributions has a χ^2 of 209 for 190 degrees of freedom, with no visible structures in the fit residuals and pull plots consistent with unit-width Gaussians.

To test the sensitivity and accuracy of the fitting procedure, the data are used to create ensembles of charge-randomized pseudoexperiments with a range of different input raw asymmetries. These confirm that the asymmetry extraction is unbiased and that the statistical uncertainty reported by the fit is consistent with the expected value ($\pm 0.15\%$). Systematic uncertainties are evaluated for a range of sources by repeating the fit under several reasonable variations and examining the change in the extracted raw asymmetry. The contribution to the systematic uncertainty on A from each source is taken as the rms of the set of fit variants with respect to the nominal measurement. The upper and lower limits of the fitting range are independently varied by up to $50 \text{ MeV}/c^2$; the bin width is varied from 2 to $10 \text{ MeV}/c^2$; an alternative method is used to determine the magnet polarity weights, based on the number of fitted signal candidates (rather than the total yield) in each configuration; the combinatorial background model is varied, either by removing the cubic term, or by adding a quadratic term; and, finally, the inflection point of the hyperbolic tangent function is

allowed to vary in the fit, rather than being fixed from simulation. The dominant systematic effect comes from varying the fitting range ($\pm 0.017\%$). Variations on the choice of bin width and fitting model contribute $\pm 0.005\%$ each, and the polarity weighting method gives an uncertainty an order of magnitude smaller. The final systematic uncertainty on A , given by summing the individual contributions in quadrature, is $\pm 0.018\%$, much smaller than the statistical uncertainty.

The detector asymmetry has one term for each final-state particle, including the muon requirement, $A_{\text{det}} = 2a^\pi + \rho \cdot a^\mu - a^K$, where a^X is the reconstruction asymmetry for particle species X . The factor of 2 accounts for the two pions in the final state, and the sign of each term reflects the charge with respect to the D meson. The muon asymmetry coefficient ρ is the charge correlation between the muon and D meson, necessary because no explicit charge requirements are enforced in this analysis. This is extracted from the data, through separate fits of the two cases $q(\mu) \cdot q(D) = \pm 1$, yielding $\rho = -0.435 \pm 0.004$, with consistent values found when analyzing D^+ and D^- samples separately. Each of the three asymmetries a^X is extracted from dedicated independent channels and determined in appropriate kinematic bins to allow them to be applied to the signal channel by a weighted average over all bins. These input asymmetries have already been determined, documented [10] and used in several previous D0 publications [8,10,19,20].

The kaon asymmetry is at least 20 times larger than all other detector effects. It arises from the larger K^- cross section with detector material than for K^+ , leading to a higher K^+ reconstruction efficiency. This asymmetry is extracted from $K^{*0} \rightarrow K^- \pi^+$ decays, in bins of absolute kaon pseudorapidity $|\eta(K)|$ and momentum $p(K)$ [10]. Applying these to the signal sample gives a total kaon asymmetry of $a^K = (1.06 \pm 0.04 \pm 0.05)\%$. The first uncertainty is statistical, from the finite K^{*0} sample size; the second uncertainty is systematic, based on variations of the K^{*0} fitting method. The pion asymmetry is investigated using $K_S^0 \rightarrow \pi^+ \pi^-$ and $K^{*+} \rightarrow K_S^0 \pi^+$ decays [10]. No indication of any asymmetry is observed, and we assign a systematic uncertainty of $\pm 0.05\%$ to account for the limited precision of this measurement.

The muon asymmetry is extracted from $J/\psi \rightarrow \mu^+ \mu^-$ decays, in bins of absolute muon pseudorapidity $|\eta(\mu)|$ and transverse momentum $p_T(\mu)$ [10]. After convoluting the kinematically binned muon asymmetry with the corresponding signal distributions, and multiplying by the charge correlation, the final correction is $\rho \cdot a^\mu = (-0.045 \pm 0.011 \pm 0.004)\%$. The systematic uncertainty includes variations to the J/ψ fitting procedure and to the kinematic binning scheme. The overall detector asymmetry is then $A_{\text{det}} = [-1.11 \pm 0.04(\text{stat}) \pm 0.07(\text{syst})]\%$, where statistical and systematic uncertainties from each source have been separately added in quadrature.

After correcting for detector asymmetries, we consider the asymmetry A_{phys} arising from different rates of D^+ and D^- production. We assume that the direct production of D^\pm mesons from $c\bar{c}$ (and B mesons from $b\bar{b}$) is charge symmetric. We also assume that there is negligible CP violation in the decays of B mesons into final states containing a D^\pm meson. We allow a contribution to the D^\pm production asymmetry from CP violation in the mixing of neutral $B_{d(s)}^0$ mesons, quantified by the mixing asymmetry parameters $a_{\text{sl}}^{d(s)}$ which are taken to be the current world averages $a_{\text{sl}}^d = (-0.09 \pm 0.21)\%$ and $a_{\text{sl}}^s = (-0.77 \pm 0.42)\%$ [4].

To determine the fraction of D mesons in our sample that originate from such decays, we use MC simulation, passed through the full data reconstruction and reweighted to match the data in five important variables: the muon multiplicity, $p_T(\mu)$, $|\eta(\mu)|$, $q(\mu) \cdot q(D)$, and the separation of the muon and D meson along the beam direction (at their point of closest approach in the transverse plane). The simulation is of $D^\pm \rightarrow K^\mp \pi^\pm \pi^\pm$ decays with the muon requirement only placed during simulation of the trigger and off-line event selection, to ensure a representative mixture of muons from the initial hard scatter, from decays of heavy-flavor hadrons, and from decays of charged kaons and pions. A fraction $(52.3 \pm 0.3)\%$ of D mesons is found to originate from the decays of B^0 mesons, but only $(12.1 \pm 0.2)\%$ from B^0 mesons that oscillated into their antiparticle prior to decay. For B_s^0 mesons, the corresponding fractions are $(2.7 \pm 0.1)\%$ total and $(1.33 \pm 0.06)\%$ oscillated. Multiplying by the respective mixing asymmetries, the contributions to A_{phys} are $(-0.010 \pm 0.023)\%$ from B^0 and $(-0.004 \pm 0.002)\%$ from B_s^0 mesons. The uncertainties are dominated by the $a_{\text{sl}}^{d(s)}$ inputs and are taken as systematic. All other reasonable variations to the method (modified reweighting, adjusted lifetimes, mixing frequencies, and branching fractions) give negligible shifts with respect to the precision. Adding these contributions, we obtain $A_{\text{phys}} = (-0.014 \pm 0.023)\%$. Of the remaining D mesons, $(35.9 \pm 0.3)\%$ arise from direct $c\bar{c}$ hadronization, $(9.0 \pm 0.2)\%$ are from B^\pm decay, and the remaining $(0.10 \pm 0.02)\%$ are from b baryons. For all cases, the uncertainties on the quoted fractions come from the limited statistics of the simulation.

From Eq. (2), we obtain the final measurement

$$A_{CP}(D^+ \rightarrow K^- \pi^+ \pi^+) = [-0.16 \pm 0.15(\text{stat}) \pm 0.09(\text{syst})]\%. \quad (4)$$

This result is consistent with the standard model prediction of CP conservation. In this evaluation, only the statistical uncertainty on A is included in the final statistical uncertainty on A_{CP} . All other uncertainties are taken to be systematic, since they are not directly related to the size of the signal sample. They are added in quadrature and treated

as completely uncorrelated, with contributions of $\pm 0.018\%$ from A , $\pm 0.084\%$ from A_{det} , and $\pm 0.023\%$ from A_{phys} . Because of the $p_T > 0.7 \text{ GeV}/c$ requirements on the three final-state hadrons, we observe some efficiency variation over the Dalitz plane. However, the efficiency is fairly uniform over $\approx 80\%$ of the allowed phase space, with relative changes of around $\pm 15\%$ in this region, and so should not select significantly against individual contributing amplitudes.

We perform a range of cross-checks to demonstrate the stability of the measurement by repeating the entire analysis for orthogonal subsamples of the data, divided in important variables including the LLR discriminant output, positive and negative kaon pseudorapidity, $p(K)$, $|\eta(K)|$, $q(\mu) \cdot q(D)$, and the instantaneous luminosity. In total, 19 such samples are tested, and all A_{CP} measurements are consistent with the nominal value.

In conclusion, we have measured the direct CP -violating parameter in the Cabibbo-favored decay $D^+ \rightarrow K^- \pi^+ \pi^+$, finding an asymmetry consistent with the SM prediction of zero. The precision exceeds that of the previous best measurement by a factor of 2.5 and represents an important reference measurement for future studies of CP violation in charm and bottom hadron decays. In particular, it gives experimental confirmation of the assumptions used in measurements of CP violation in D^0 and B^0 mixing and decay [10,11], which is of special importance given the anomalously large asymmetry reported in same-charge dimuons [5], and for future searches for CP violation in bottom and charm hadrons.

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